

THE DYNAMICS OF MODULATED

WAVE TRAINS

(in reaction-diffusion equations)

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* What's a 'wave train'?

- Why do they appear in families?
- Examples (cGL, FH-N)

* Modulated wave trains

- (Weak) shocks / defects ↔ 'SINKS'

* The derivation of Burgers equation as

PHASE DIFFUSION EQUATION

- Stability of the wave train
- 'Ansatz'
- Solvability

* VALIDITY

* Shocks & defects

- In Burgers eq
- In the full system ↔ group speed

WHAT'S A WAVE TRAIN?

R-d system: $U_t = D U_{xx} + F(U)$

$$U: \mathbb{R} \times \mathbb{R}^+ \rightarrow \mathbb{R}^N$$

Note: $x \in \mathbb{R}$

D = diffusion matrix

$F(U)$: reaction / nonlinearities

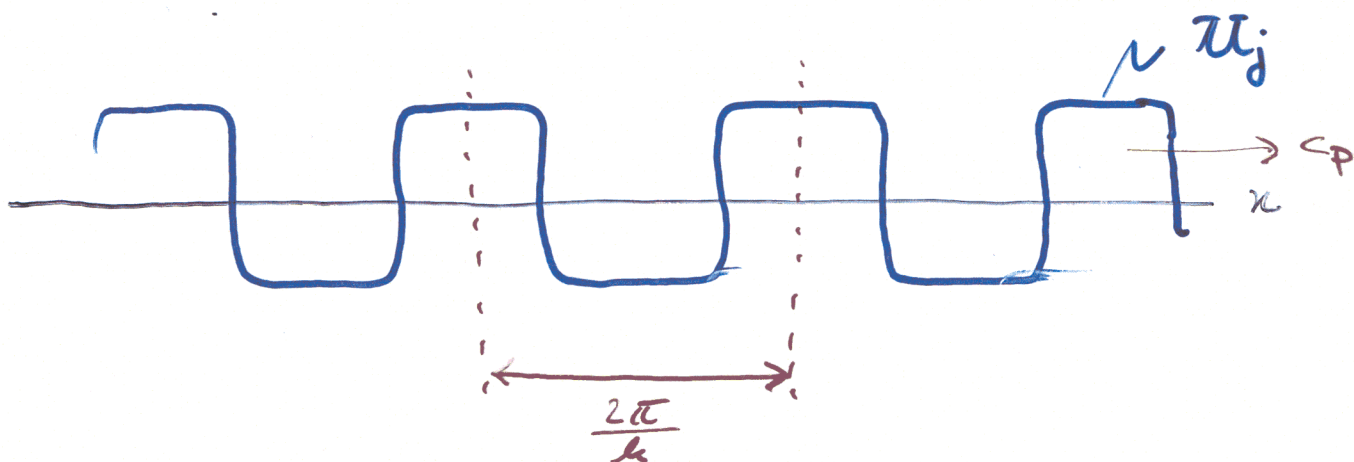
$$U(x, t) = u_0(\omega t - kx) = u_0(\phi)$$

a wave train

2π -per ω in ϕ

$$-k(x - \frac{\omega}{k}t)$$

$$c_p = \frac{\omega}{k} \quad \text{"phase speed"}$$



CLAIM: In general there is a 1-parameter family of wave trains

$$U(x,t) = u_0(\omega(k)t - kx; k)$$

WHY?

Eq for $u(\phi)$ ($\phi = \omega t - kx$)

$$\boxed{k^2 D u_{\phi\phi} - \omega u_{\phi} + F(u) = 0}$$

+ per. b.c. on $(0, 2\pi)$

a 2N-dim ODE

Let's vary k : $\frac{\partial}{\partial k}$ evaluated at u_0, ω_0, k_0

$$u_k = \frac{\partial u}{\partial k} \Big|_{k=k_0}$$

$$\omega_k = \frac{\partial \omega}{\partial k} \Big|_{k=k_0}$$

$$\mathcal{L}_0 u_k = [k_0^2 D \partial_{\phi}^2 u_0 - \omega_0 \partial_{\phi} u_0 + F'(u_0)] u_k$$

$$= \omega_k \partial_{\phi} u_0 - 2k_0 \partial_{\phi}^2 u_0$$

on $L_2(0, 2\pi)$

Operator associated to lin. stab. of $u_0(\phi)$

ASSUME: generalized kernel of $\mathcal{L}_0$ is 1-d

$\Rightarrow$ spanned by $\partial_{\phi} u_0$

$\uparrow$ translation invariance

we

Adjoint operator u_{ad}

(u_0 in general not self-adj.)

Kernel of L_{ad} spanned by u_{ad}

normalization: $\langle \partial_\phi u_0, u_{ad} \rangle_{L_2} = 1$

Now $L_0 u_k = b$ ("inhomogeneous")

$\Rightarrow$ solvability criterion: $\langle b, u_{ad} \rangle = 0$

(since: $\langle b, u_{ad} \rangle = \langle L_0 u_k, u_{ad} \rangle = \langle u_k, L_{ad} u_{ad} \rangle$)

existence $\frac{\partial u}{\partial k}$

$\Downarrow$

$\Rightarrow \langle \omega_k \partial_\phi u_0 - 2k_0 D \partial_\phi^2 u_0, u_{ad} \rangle = 0$

$\Downarrow$

$$\omega_k = 2k_0 \langle D \partial_\phi^2 u_0, u_{ad} \rangle$$

$\uparrow$

$$\boxed{\omega_k = \frac{\partial \omega}{\partial k} = c_g : \text{GROUP SPEED}}$$

Thus: for this ω_k , $L u_k = b$ can be solved
(uniquely up to transl.)

$\updownarrow$
 $\hookrightarrow \partial_\phi u_0$

With this assumption on ω there is 1-parameter fam.

Not'n $\omega = \omega_{ne}(k)$ "nonlinear dispersion"

EXAMPLE : complex Ginzburg - Landau

$$A_t = (1 + i\alpha) A_{xx} + A - (1 + i\beta) |A|^2 A$$

$$A = u_1 + i u_2 \rightarrow \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} : \mathbb{R} \times \mathbb{R}^+ \rightarrow \mathbb{R}^2$$

'near equilibrium'

α, β : parameters

$\uparrow$ 'the underlying system'

Wave trains : $A_0(x, t) = R e^{i(\omega t - kx)}$

$$\rightarrow \begin{cases} R^2 + k^2 = 1 \Rightarrow R = R(k) = \sqrt{1 - k^2} \\ \omega = \omega(k) = \beta - (\alpha - \beta) k^2 \end{cases}$$

$$\uparrow \boxed{\omega_{ne}(k) = \beta - (\alpha - \beta) k^2}$$

Note $\omega_{nl} \equiv 0$ in RGL - case ($\alpha = \beta = 0$)

Note $A_0(\omega t - kx; k)$ is stable if

$$1 + \alpha\beta > 0$$

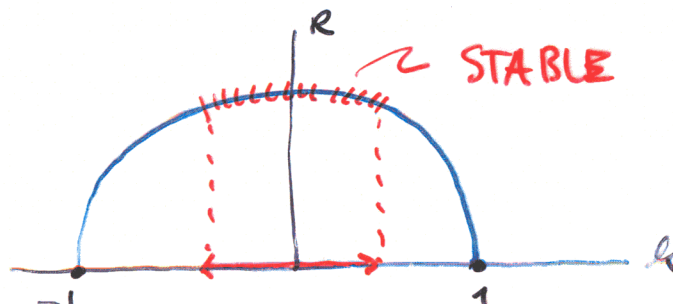
$\&$

$$k^2 < \frac{1 + \alpha\beta}{3 + \alpha\beta + 2\beta^2}$$

Benjamin - Feir

Eckhaus,

Newell, ...



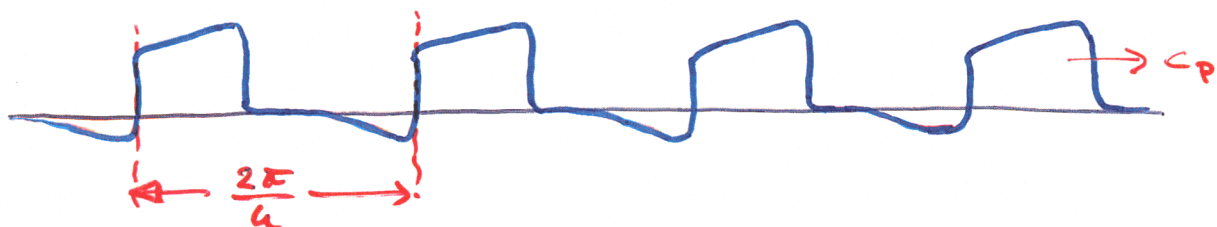
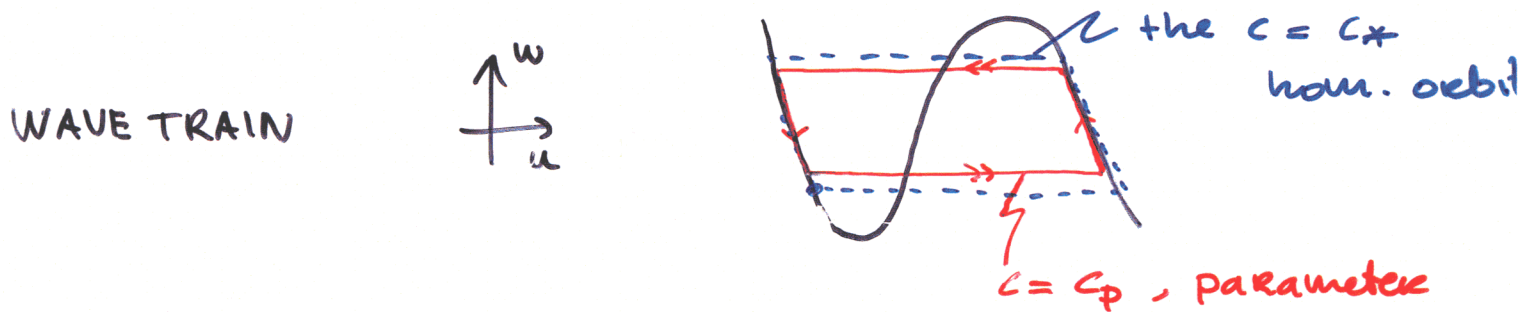
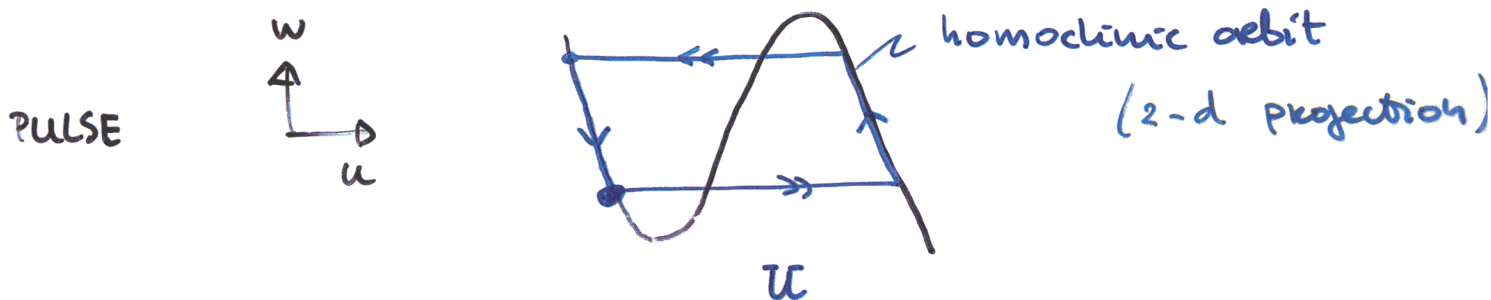
FitzHugh - Nagumo $\leftrightarrow$ far-from-equilibrium

$$\begin{cases} u_t = u_{xx} + u(1-u)(u-a) + w \\ w_t = \varepsilon(u - b w) \end{cases}$$

$\rightarrow \exists$ traveling LOCALIZED pulse (& it is stable [Jones])

$\updownarrow$
well-defined c_*

[Hastings] For c near c_* there exist a family of traveling wave trains



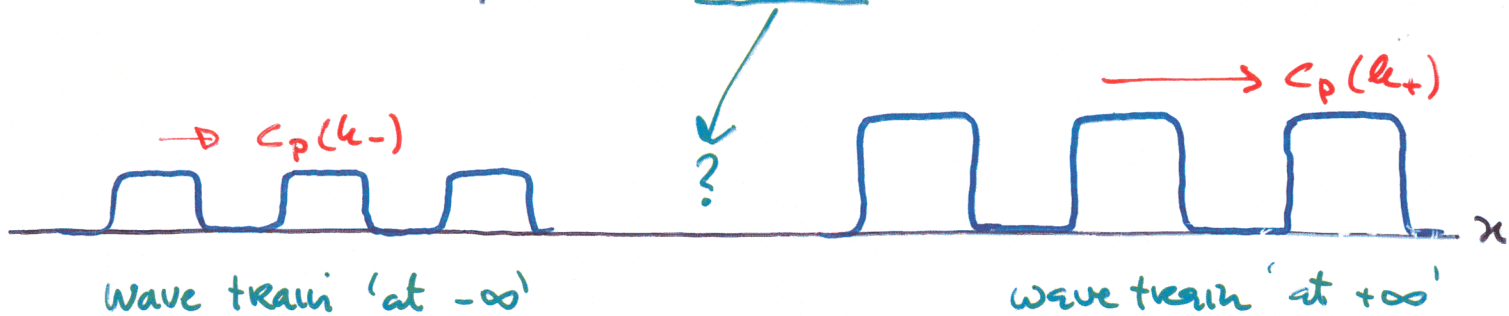
stability: Eszter & Jandhe

MODULATED WAVE TRAINS



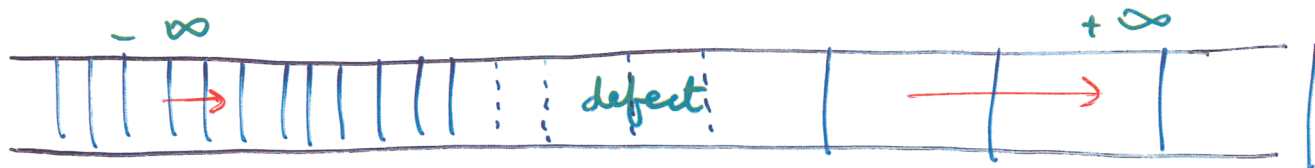
DYNAMICS (of full PDE) IN, OR NEAR, THE FAMILY OF TRAVELING WAVE TRAINS

For instance: Q / Can 'defects' exist (& be stable)?



From above (semi-2-D)

"heteroclinic orbit between wave train"



ANSWER: Yes, but it depends on the GROUP SPEEDS

i.e. 'Defects' or 'weak shocks' exist & are stable

$$\text{if } c_- = \frac{\partial \omega_{nl}}{\partial k}(-\infty) > c_+ = \frac{\partial \omega_{nl}}{\partial k}(+\infty)$$



SINKS

&
defect travels with
speed $c^* \in (c_-, c_+)$

"information is transported into the defect"

THE DERIVATION OF BURGERS EQUATION

THE DYNAMICS OF MODULATED WAVE TRAINS

I STABILITY OF A WAVE TRAIN

$$\mathcal{U}_t = D \mathcal{U}_{xx} + F(\mathcal{U})$$

pick a $k_0, \omega_0 = \omega_{nl}(k_0)$

in $\phi = \omega_0 t - k_0 x$ frame

(that travels along with speed c_p)

$$\mathcal{U}_t = k_0^2 D \mathcal{U}_{\phi\phi} - \omega_0 \mathcal{U}_{\phi} + F(\mathcal{U})$$

linearize around $u_0(\phi; k_0) : \mathcal{U} = u_0 + u$

$$u_t = k_0^2 D u_{\phi\phi} - \omega_0 u_{\phi} + \underbrace{F'(u_0(\phi))}_{2\pi\text{-periodic in } \phi} u$$

2π -periodic in ϕ



FLOQUET THEORY

$$u \leftrightarrow e^{-\frac{\gamma}{k_0} \phi} v(\phi) e^{\lambda t}$$

$\underbrace{v(\phi)}_{2\pi\text{-periodic}}$

$\mathcal{L}_{lin}(\gamma) v$
||

$\gamma \in i\mathbb{R} : \text{Floquet exponent}$

$$\mathcal{L}_{\gamma} v = \left[k_0^2 D \left(\partial_{\phi} - \frac{\gamma}{k_0} \right)^2 - \omega_0 \left(\partial_{\phi} - \frac{\gamma}{k_0} \right) + F'(u_0(\phi)) \right] v$$

NOTE : $\mathcal{L}_0 = \mathcal{L}_0 \leftrightarrow \text{existence pb.}$

Assumption on u_0



(gen) kernel is spanned
by $\partial_\phi u_0$

of eigenvalues $\lambda = \lambda_{\text{lin}}(\nu)$

near $\lambda = 0$



same argument as in ex. th. 

$\lambda = \lambda_{\text{lin}}(\nu)$ ($\nu \in i\mathbb{R}$): the LINEAR dispersion relation

Derivative with respect to ν evaluated at $\nu = 0$

(note: $\nu|_{\nu=0} = \partial_\phi u_0$)

$$\mathcal{L}_0 \nu_\nu = \left(\lambda'_{\text{lin}}(0) - \frac{\omega_0}{k_0} \right) \partial_\phi u_0 + 2k_0 D \partial_\phi \phi u_0$$

$$\left(\text{Ex } \mathcal{L}_0 u_k = \omega'_{nl}(k_0) \partial_\phi u_0 - 2k_0 D \partial_\phi \phi u_0 \right)$$

SOLVABILITY: $\langle \text{inhomogeneous } u_{\text{ad}} \rangle_{L_2} = 0$

$$\lambda'_{\text{lin}}(0) - \frac{\omega_0}{k_0} = -2k_0 \langle D \partial_\phi \phi u_0, u_{\text{ad}} \rangle$$

$\parallel$

$c_p(k_0)$

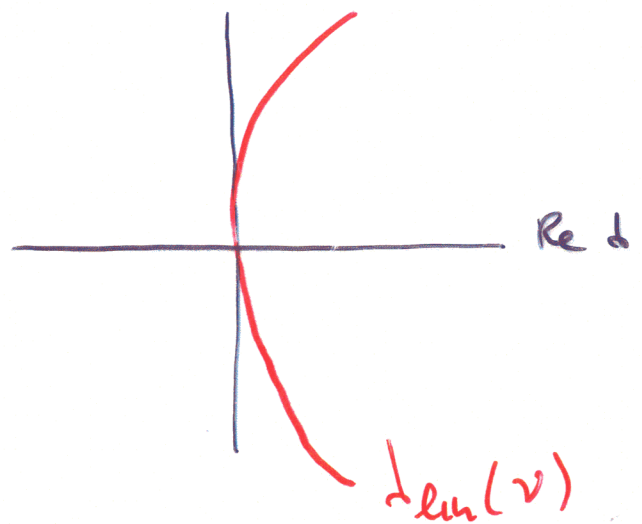
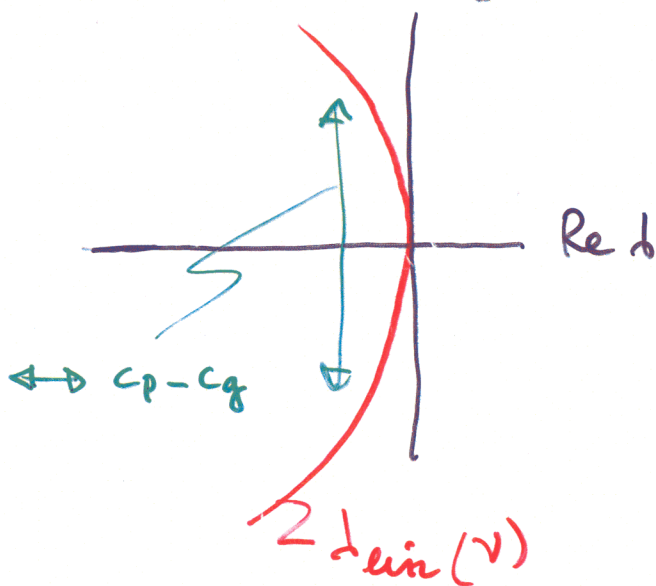
$\parallel$

$-\omega'_{nl}(k_0)$

$$\Rightarrow \lambda'_{\text{lin}}(0) = c_p(k_0) - c_g(k_0)$$

& similar relations between

$\lambda''_{\text{lin}}(0)$
 $\parallel$

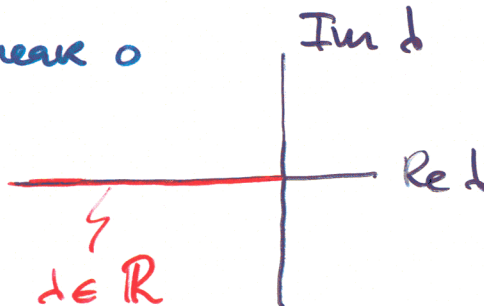


STABLE : $d''_{lin}(0) > 0$

↑
↓
(at least "locally")

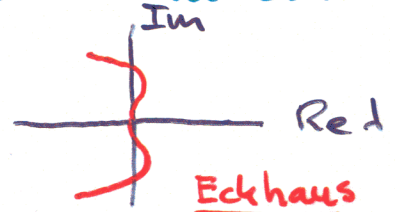
↑
↓
 d 's near 0

NOTE : $c_p = c_g$



UNSTABLE : $d''_{lin}(0) < 0$

NOTE : destabilization



'dissipative' / STABLE

Assumption

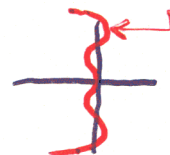
$\omega''_{nl}(k_0) \neq 0$ & $d''_{lin}(0) > 0$

↑
↓
Ex: cGL, not rGL (later)

Assumption

For ANY $v \in i\mathbb{R}$, $v \neq ik_0/2$, we have

$\text{Re } d < 0$ ("global stability")



NOTE : 'In practice' very hard ↔ van der Ploeg

↑

for a given model

(in Greek-Meinhard)

II: MODULATED WAVES: THE ANSATZ

Wave train: $\mathcal{U} = u_0(\omega_{nl}(k)t - kx - \Phi; k)$, k par.

Φ
phase, a priori a free constant

Assume Φ is allowed to evolve slowly in x & t ,
so that $\mathcal{U}(x,t)$ remains in/near the fam. of wave trs.

Introduce $\delta \ll 1 : k \rightarrow k + \delta$
 $\uparrow$ 'ARTIFICIAL'

in a small nbhd
of original k

$\Rightarrow u_0(\omega(k)t - kx; k) \rightarrow$ a nearby wave tra

$$u_0(\omega(k+\delta)t - (k+\delta)x; k+\delta)$$

$$= u_0([\omega(k)t + \delta\omega'(k) + O(\delta^2)]t - [k+\delta]x; \cdot)$$

$$= u_0(\underbrace{[\omega(k)t - kx]}_{\delta=0} - \underbrace{\delta[x - c_g(k)t]}_{X \text{ slow space}} + \underbrace{\delta^2 t [\dots]}_{T \text{ slow time}})$$

$$\Phi = \Phi(X, T)$$

$$= \Phi(\delta(x - c_g t), \delta^2 t)$$

"natural (slow) x & t scales of phase"

Note: $\Phi(X, T)$ induces change in k ("vice versa")

$$\begin{aligned}
 kx + \Phi &= kx + \Phi_0 + X\Phi_X + O(\delta^2) \\
 &= \underbrace{(k + \delta\Phi_X)}_{\text{change in } k} x + \Phi_0 - \underbrace{c_g t \delta}_{\substack{\text{induced} \\ \text{related change} \\ \text{in } \omega_{nl}(k)}} + O(\delta^2)
 \end{aligned}$$

$\downarrow X^2, T$

$\Rightarrow$ ANSATZ

$$X = \delta(x - c_g t), T = \delta^2 t$$

$$\begin{aligned}
 \mathcal{U}(x, t) &= u_0(\omega_{nl}(k)t - kx - \Phi(X, T); \underbrace{k + \delta\Phi_X}_{\text{change in } k}) \\
 &\quad + \underbrace{\delta^2 u_2(\omega_{nl}t - kx; X, T)}_{\substack{\text{rest} \\ \text{"remainder"}}}
 \end{aligned}$$

SUBSTITUTE THIS ANSATZ INTO $\mathcal{U}_t = D\mathcal{U}_{xx} + F(\mathcal{U})$

$$\begin{aligned}
 \text{Note: } \frac{\partial}{\partial t} \mathcal{U} &= \frac{\partial}{\partial t} [u_0(\omega - \Phi(X, T); k + \delta\Phi_X(X, T)) + \delta^2 u_2(\phi, X, T)] \\
 &= \partial_\phi u_0 [\omega + \delta c_g \Phi_X - \delta^2 \Phi_T] + \partial_k u_0 [-\delta^2 c_g \Phi_{XX}] \\
 &\quad + \delta^2 \partial_\phi u_2 [\omega] \\
 &\quad + O(\delta^3)
 \end{aligned}$$

$\rightarrow$ ETCETERA $\leftarrow$

III THE SOLVABILITY CONDITION

$$0 = -u_t + D u_{xx} + F(u)$$

$$= \dots$$

Existence pb.

$$= -\omega \partial_\phi u_0 + k^2 D \partial_\phi^2 u_0 + F(u_0)$$

Exist. pb ($\exists$ fam.)

$$+ \delta [L_0 \partial_k u_0 + 2k D \partial_\phi^2 u_0 - c_g \partial_\phi u_0]$$

"correct speed"

$$+ \delta^2 [L_0 u_2 + \Phi_T [\partial_\phi u_0] +$$

$$\Phi_{xx} [\dots] +$$

$$(\Phi_x)^2 [\dots]] + O(\delta^3)$$

$$\Rightarrow \boxed{L_0 u_2 = -\Phi_T [\dots] - \Phi_{xx} [\dots] - (\Phi_x)^2 [\dots]}$$

L_0 has nontrivial kernel $\partial_\phi u_0 \iff$

L_{ad} has kernel spanned by u_{ad}

$$\& \langle \partial_\phi u_0, u_{ad} \rangle_{L_2} = 1$$

$$L_0 u_2 = b \Rightarrow \text{only solvable if}$$

$$\langle b, u_{ad} \rangle_{L_2} = 0$$

$$\Phi_T = \langle -[\dots], u_{ad} \rangle \Phi_{xx} + \langle -[\dots], u_{ad} \rangle (\Phi_x)^2$$



identities for $d'_{en}(0), d''_{en}(0)$
 $\& \omega'_{nl}(k), \omega''_{nl}(k)$

$$\Phi_T = \frac{1}{2} d''_{en}(0) \Phi_{xx} - \frac{1}{2} \omega''_{nl}(k) (\Phi_x)^2$$

PHASE DIFFUSION EQUATION

Def : $q(x, T) = \Phi_x(x, T)$

"wave number modulation"

$$\Rightarrow q_T = \frac{1}{2} d''_{en}(0) q_{xx} - \omega''_{nl}(k) q q_x$$

BURGERS EQUATION (!)

Recall : we assumed

$$d''_{en}(0) > 0$$



$$\omega''_{nl}(k) \neq 0$$



STABILITY

Th (For each $G > 0, L > 0$ & $R \in [0, 2) \exists \delta_0, G_0 > 0$
st the following holds:)

Let $q(X, T)$ be a solution of

$$q_T = \frac{1}{2} d''(0) q_{XX} - \omega''(k) q q_X$$

s.t. $\sup_{[0, T_0]} \|q(\cdot, T)\|_X \leq G$

define $\Phi(X, T) = \int_0^X q(Y, T) dY$ &

$$u_{app.}(x, t) = u_0(\omega t - kx - \Phi(X, T); k + \delta q(X, T))$$

Then there exists a sol. u of $u_t = D u_{xx} + F(u)$

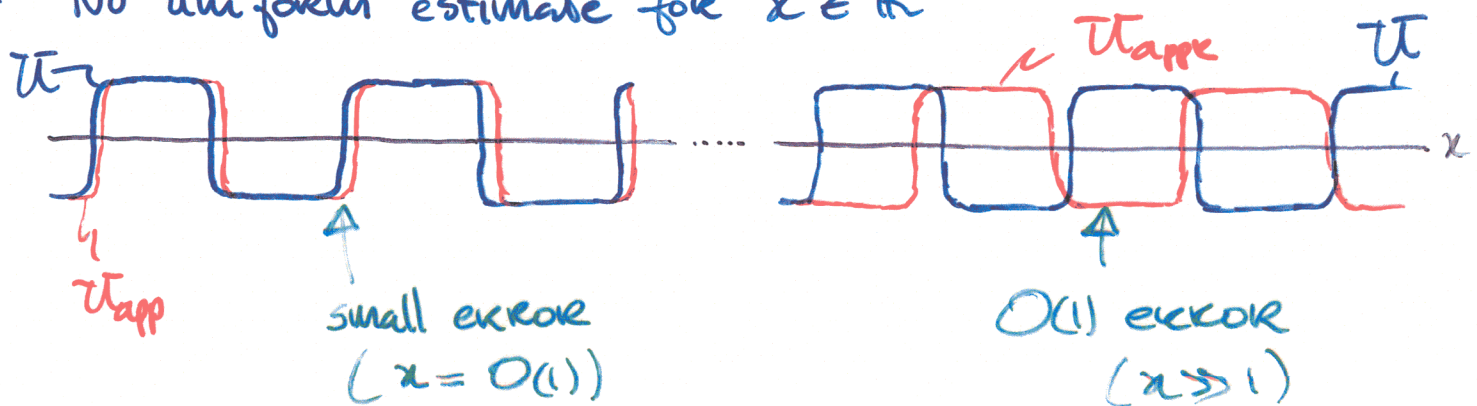
s.t. $\sup_{t \in [0, T_0/\delta^2], |x| < L/\delta^R} |u(x, t) - u_{app.}(x, t)| \leq G_0 \delta^{3-R}$

for each $\delta \in (0, \delta_0)$

Note * Validity on t -scale Burgers & beyond its x -scale
($R = 3/2$)

* " Gronwall" $\rightarrow O(\frac{1}{\delta})$ t -scale, useless

* No uniform estimate for $x \in \mathbb{R}$



Burgers $q_T = \frac{1}{2} d'' q_{xx} - \omega'' q q_x$

Stationary fronts ("we're already traveling with c_f ")
(*)

→ if K s.t. $\text{sign } K = \text{sign } \omega''$ then

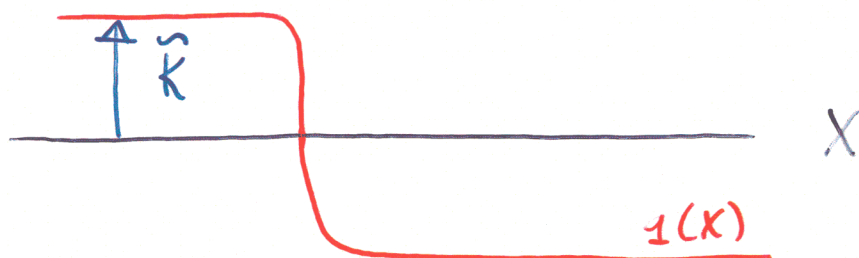
$$q(x) = \alpha \tanh \beta x$$

with $\alpha = -\text{sign } \omega'' \sqrt{\frac{K}{\omega''}}$

$$\beta = \frac{1}{d''} \sqrt{K \omega''}$$

def $\tilde{K} > 0$

$\omega'' > 0$



$\omega'' < 0$



All fronts ($\exists$ families) are (non-) stable

$\Updownarrow$
 $\tilde{K} > 0$

(*) Traveling front $\iff$ stationary front

for Burgers eq 'centered' around $\tilde{k}_0 = k_0 + O(\epsilon)$

slightly different c_f

$$\omega'' > 0 \Rightarrow q \xrightarrow{X \rightarrow \pm\infty} \mp \tilde{K}, \pm \tilde{K}$$

$$\Rightarrow \Phi \xrightarrow{X \rightarrow \pm\infty} \mp \tilde{K}X, \pm \tilde{K}X$$

(+ $\Phi_0 t$..)

$$\Rightarrow \omega'' > 0 \Rightarrow u \xrightarrow{X \rightarrow \pm\infty} u_0(\omega t - kx \pm \tilde{K}X; k \pm \delta \tilde{K})$$

$$\omega'' < 0 \Rightarrow u \xrightarrow{X \rightarrow \pm\infty} u_0(\omega t - kx \mp \tilde{K}X; k \pm \delta \tilde{K})$$

$\delta(x - c_g t)$

$\Rightarrow \omega'' > 0$: shock/defect connects wave train

"at $-\infty$ ": $u_0(\underbrace{[\omega \pm \delta c_g \tilde{K}]t - [k \pm \delta \tilde{K}]x}_{\parallel \text{ (approx)}}; k \pm \delta \tilde{K})$

$$u_0(\omega(k \pm \delta \tilde{K})t - (k \pm \delta \tilde{K})x; k \pm \delta \tilde{K})$$

to "at $+\infty$ ": $u_0(\omega(k \mp \delta \tilde{K})t - (k \mp \delta \tilde{K})x; k \mp \delta \tilde{K})$

THUS $\omega'' > 0 \Rightarrow c_g = \omega'$ increasing
decreasing as fct. of k

$c_g(-\infty) > c_g(k) > c_g(+\infty)$

$$\omega'' > 0 : \begin{cases} k + \delta \tilde{K} \\ k - \delta \tilde{K} \end{cases} \xrightarrow{\text{red}} k - \delta \tilde{K} \xrightarrow{\text{green}} k + \delta \tilde{K}$$

$-\infty \qquad \qquad \qquad +\infty$

$\tilde{K}$ par $\tilde{K} > 0$

$\Rightarrow$ Burgers equation gives approximation (on finite t & x-domains!)
 of SHOCK / DEFECTS with $\uparrow$
 $!!$

$$c_{\text{group at } -\infty} > c_{\text{defect}} > c_{\text{group at } +\infty}$$

$\updownarrow$
SINKS

(Rankine-Hugoniot)

$$k_{\pm} = k \mp \delta \tilde{K}$$

$$k_{\pm} = k \pm \delta \tilde{K}$$

Note : $c_g(k) = c_{\text{defect}} = \frac{\omega(k_{+}) - \omega(k_{-})}{k_{+} - k_{-}} = \dots$

where $k_{\pm} = k|_{at \pm \infty}$

CAN WE PROVE THE EXISTENCE & STABILITY
 OF SUCH SHOCKS / DEFECTS FOR THE FULL
EQUATION (& FOR $T \geq 0, x \in \mathbb{R}$) ?

UNBOUNDED!

NOTE : BURGERS GIVES DYNAMICS OF ALL POSSIBLE
 PHASE EVOLUTIONS (bounded T & X)

NOW : EX & STAB OF A SPECIAL KIND OF BEHAVIOR

TRAVELING TIME - PERIODIC SOLUTIONS

$$U(x, t) = u_* (\underbrace{x - c_* t}_{\text{trav. coord } \xi}, \underbrace{\omega_* t}_{\tau}, 2\pi - \text{per.})$$

u_* must be bi-asymptotic to wave trains

$$u_*(x - c_* t, \omega_* t) \xrightarrow{x \rightarrow \pm \infty} u_0(\omega_{\pm} t - k_{\pm} x; k_{\pm})$$

More precisely :

$$\left\| u_* \left(\underbrace{\xi}_{\substack{\uparrow \\ \tau = \omega_* t}}, \cdot \right) - u_0 \left(\frac{\omega_{\pm} - c_* k_{\pm}}{\omega_*} \cdot - k_{\pm} \xi; k_{\pm} \right) \right\|_{H^1(0, 2\pi)} \xrightarrow{\xi \rightarrow \pm \infty} 0$$

$x = \xi + c_* t$

⇒ JUMP CONDITIONS :

u_0 must have frequency 1 in τ !

$$\frac{\omega_+ - c_* k_+}{\omega_*} = 1 = \frac{\omega_- - c_* k_-}{\omega_*}$$

⇒

$$c_* = \frac{\omega_+ - \omega_-}{k_+ - k_-} \quad (\leftrightarrow \text{Burgers})$$

$$\omega_* = \omega_{\pm} - c_* k_{\pm}$$

Th (Assumptions on λ'' , ω'' , stab.)

For all wave numbers k_- & k_+ close to k_0

s.t. $\boxed{c_g^- = \omega'(k_-) > c_g^+ = \omega'(k_+)}$

there exists a $\left\{ \begin{array}{l} \text{weak viscous shock} \\ \text{defect} \\ \text{soliton} \end{array} \right\}$ solution

$$U(x, t) = u_*(x - c_* t, \omega_* t) \quad \text{with}$$

$$c_* = \frac{\omega(k_+) - \omega(k_-)}{k_+ - k_-} \in (c_g^+, c_g^-)$$

$$\omega_* = \omega(k_+) - k_+ c_*$$

u_* is unique, up to translations in x & t

MOREOVER

Th $u_*(x - c_* t, \omega_* t)$ is nonlinearly,
asymptotically stable w.r.t. perturbations
in X ($\hookrightarrow$ locally H^1)

Proof: SPATIAL DYNAMICS & CENTER MANIFOLDS



$\uparrow$ Kirchgässner

$$\begin{cases} U_3 = V \end{cases}$$

$$\begin{cases} V_3 = -D^{-1}[-\omega_* \partial_t U + c_* V + F(U)] \end{cases}$$

$$\Leftrightarrow \vec{U}_3 = \mathcal{L}(U) + N(U)$$

$\updownarrow$
ill-posed dyn. syst

APPLICATIONS

* COMPLEX GINZBURG-LANDAU

$$A_t = (1+i\alpha) A_{xx} + A - (1+i\beta) |A|^2 A$$

$$\& A_0(\omega t - kx; k) = R(k) e^{i(kx - \omega t)}$$

$$\text{Consider } k_0 = 0 \rightarrow A_0 = 1 e^{-i\omega_0 t}$$

→ Burgers :

$$\boxed{q_T = (1+\alpha\beta) q_{xx} + 2(\beta-\alpha) q q_x} + O(\delta)$$

$\updownarrow$
 $d''_{\text{en}}(0) > 0$
 $\updownarrow$
 $\boxed{1+\alpha\beta > 0}$

$\uparrow$
 $\omega''_{nl}(k_0=0) = 0$
 $\updownarrow$
 $\& \alpha = \beta = 0$
 $\updownarrow$
real GL.

Literature : in essence already in Howard & Kopell ('77)
 (no validity) & Mather & Pomeau ('79)

MOST LITERATURE FOCUSES ON $1+\alpha\beta \leq 0$ (& $\alpha = \beta = 1$)

Ex. analysis in ODE!

the transition to instability

'Rigorous shocks' : $\alpha = \beta = 0$ Briconmont & Kupiainen ('91)
 $\alpha = 0$ Kapitula ('91)

→ near the Eckhaus instability
↓

Note

$$1 + \alpha\beta \lesssim 0 \Rightarrow$$

* $\beta \neq \alpha$ ($\Leftrightarrow \omega'' \neq 0$) Burgers transforms into
(rescale X)

$$q_T = (\gamma_1 q_{xxx} + 2(\beta - \alpha)qq_x) + \delta [\gamma_2 q_{xxxx} + \gamma_3 q_{xx} + \gamma_4 (q^2)_{xx}]$$

'a perturbed KdV - eq'

↑
different KdV

* $\omega'' = 0$ (as in real GL)

Beknoob, van Haelen
(88) (94)

$$q_T = \gamma_2 q_{xxxx} + \gamma_3 q_{xx} + \gamma_4 (q^2)_{xx}$$

(rescale T)

Kuramoto - Sivashinsky eq

Kramer & Zimmerman '85

↑
a more general
setting than RGL

= FORMAL ANALYSIS -



VALIDITY ALONG THE LINES OF
OUR APPROACH TO BURGERS

OTHER APPLICATIONS

* Shocks & defects for FitzHugh-Nagumo &

other well-studied systems
(Gray-Scott, Gierer-M.)



* Systems near equilibrium



Mielke, Schneider, Coulet & Eckmann, ...

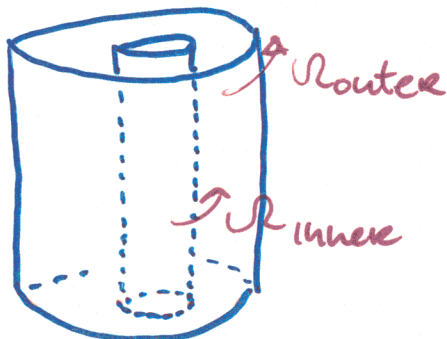
Bif. described by cgl



[DSSS]

Burgers & shocks in cgl

→ Shocks & defects in Taylor-Couette



gl bif to rings
(RGL)

or spirals
(cgl)

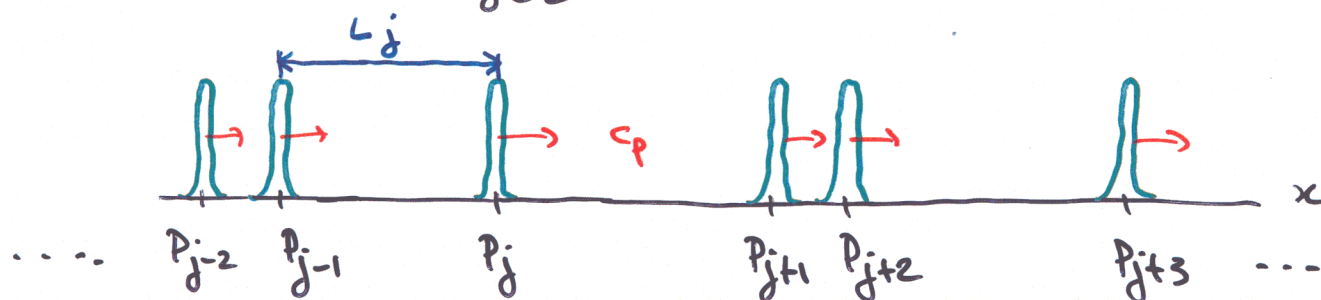
FINALLY : PULSE INTERACTIONS & LATTICE MODELS

$\updownarrow$
 the long wave length limit

Assume: $U(x,t) = h(x - c_p t)$

$\uparrow$
stable localized traveling pulse

Set: $U(x,t) \approx \sum_{j \in \mathbb{Z}} h(x - c_p t + P_j(t))$



some
cond.

$$\dot{P}_j = a e^{-b(P_j - P_{j-1})}$$

"lattice equation"

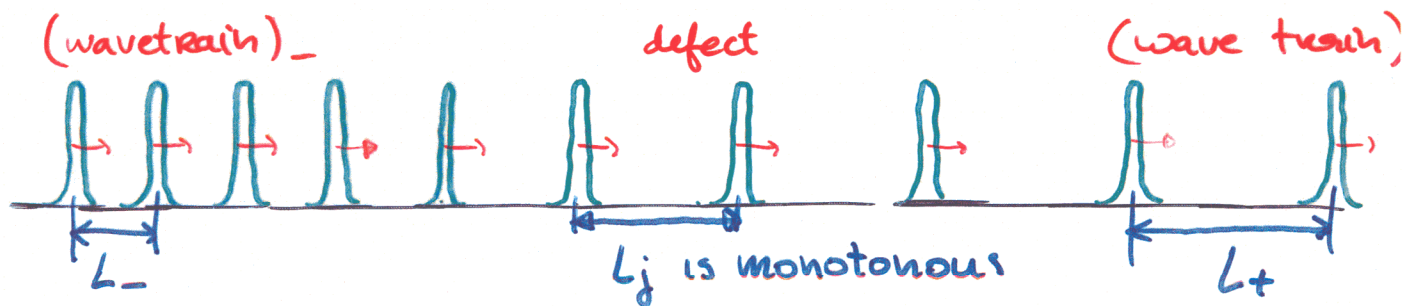
Ec, Sandstede,

Th

(Assumptions)

For any $L_+ > L_- > 0$ there exists a shock / defect solution with $L_j \rightarrow L_{\pm}$ as $j \rightarrow \pm \infty$ ($\& NOT \nexists L_+ < L_- !$)

c_g condition, Note $L_+ - L_-$ NOT SMALL



CONCLUSION / RECAPITULATION

(Slow) dynamics of modulated wave trains

"generically" captured by BURGERS eq

= at finite (& LONG) spatial & temporal scales =

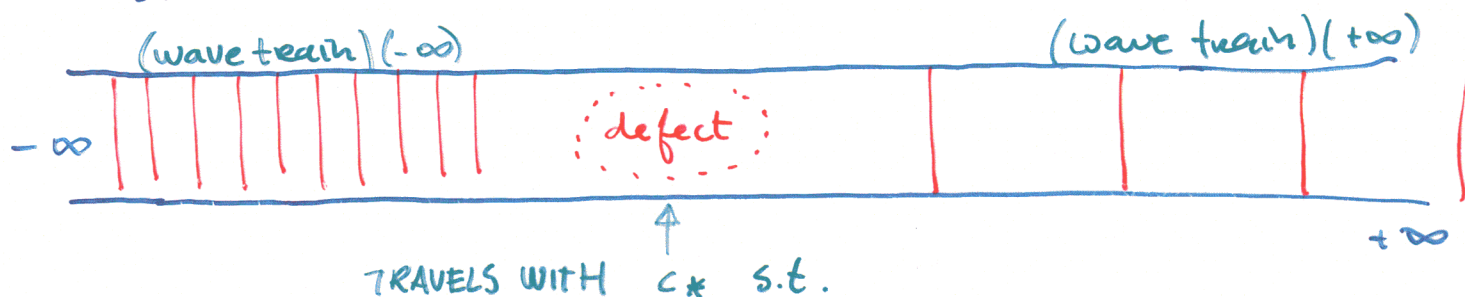
&

Existence and stability of DEFECTS OF SINK-TYPE

(for $x \in \mathbb{R}$, $t \in \mathbb{R}^+$) in full equation

(approximated by the (viscous)
shock waves of Burgers approx.)

SINKS



$$c_{\text{group}}(+\infty) < c_* < c_{\text{group}}(-\infty)$$