







Accurate Quantum State Estimation

via

"Keeping the Experimentalist Honest"

quant-ph/0603116

Robin Blume-Kohout (Caltech-IQI)

with Patrick Hayden (McGill)

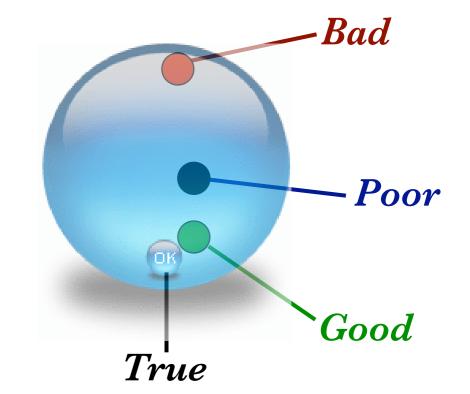
and Karan Malhotra (IIT-Kanpur)

What is State Estimation?

- **Goal**: Characterize a source of quantum systems.
- You measure N identical copies.
- Measurement record:

$$\mathcal{M} = \left\{\hat{E}_1, \hat{E}_2, \dots \hat{E}_N
ight\}$$

- Then report your best guess for ρ .
- Key Application: Verifying quantum hardware for fault tolerant quantum computation.
 - => probabilities in $[10^{-3}...10^{-5}]$ are important!



What's the Problem?

Current technology: Maximum Likelihood Estimation

$$\mathcal{M} = \left\{\hat{E}_1, \hat{E}_2, \dots \hat{E}_N
ight\} \, \longrightarrow \mathcal{L}(
ho) \equiv p(\mathcal{M}|
ho) \longrightarrow
ho_{ ext{ iny MLE}}$$

- $\rho_{\rm MLE}$ does not honestly represent the experimentalist's knowledge about the observed ensemble.
- Why? $ho_{
 m MLE}$ typically has zero eigenvalues:

$$oldsymbol{\lambda}_i = ra{\phi_i}
ho_{ ext{MLE}} \ket{\phi_i} = 0$$

Zero eigenvalue = zero probability

absolute certainty

doesn't admit error bars!

Yes, it's a problem

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Measurement of qubits

Daniel F. V. James, 1,* Paul G. Kwiat, 2,3 William J. Munro, 4,5 and Andrew G. White 2,4

1 Theoretical Division T-4, Los Alamos National Laboratory, Los Alamos, New Mexico 87545

2 Physics Division P-23, Los Alamos National Laboratory, Los Alamos, New Mexico 87545

3 Department of Physics, University of Illinois, Urbana-Champaign, Illinois 61801

4 Department of Physics, University of Queensland, Brisbane, Queensland 4072, Australia

5 Hewlett-Packard Laboratories, Filton Road, Stoke Gifford, Bristol BS34 8QZ, United Kingdom

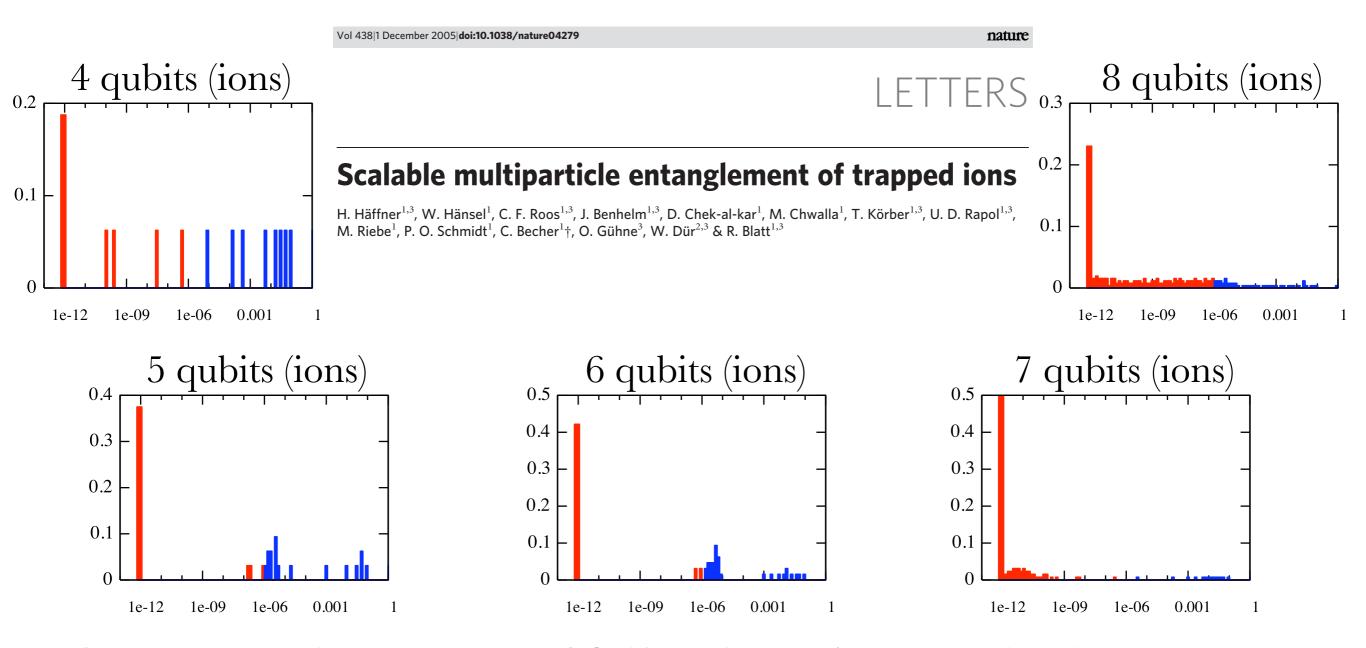
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$$\hat{\rho} = \begin{pmatrix} 0.5069 & -0.0239 + i0.0106 & -0.0412 - i0.0221 & 0.4833 + i0.0329 \\ -0.0239 - i0.0106 & 0.0048 & 0.0023 + i0.0019 & -0.0296 - i0.0077 \\ -0.0412 + i0.0221 & 0.0023 - i0.0019 & 0.0045 & -0.0425 + i0.0192 \\ 0.4833 - i0.0329 & -0.0296 + i0.0077 & -0.0425 - i0.0192 & 0.4839 \end{pmatrix}$$

This matrix is illustrated in Fig. 3 (right). In this case, the matrix has eigenvalues 0.986 022, 0.013 977 7, 0, and 0; and $\text{Tr}\{\hat{\rho}^2\}=0.972\,435$, indicating that, while the linear reconstruction gave a nonphysical density matrix, the maximum likelihood reconstruction gives a legitimate density matrix.

1. D. F. V. James et. al., "Measurement of Qubits," Phys. Rev. A, 64:052312 (2001)

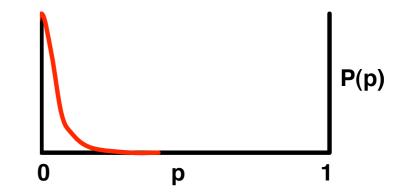
Yes, it's a problem

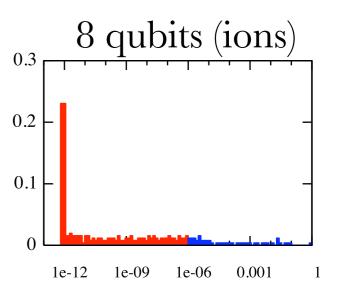


- 1. D. F. V. James et. al., "Measurement of Qubits," *Phys. Rev. A*, 64:052312 (**2001**)
- 2. Häffner et. al., "Scalable multiparticle entanglement of trapped ions," *Nature*, 438:643-6 (2005)

Error Bars?

- "But aren't there error bars in there?"
- Not always.
- What does $\frac{p=0\pm0.1}{p=0.05\pm0.05}$ mean?





States as Predictions

- Quantum states are like probability distributions: they predict the outcome of future measurements.
- Lets define a metric $f(\rho : \sigma)$ that measures how well σ predicts measurements on ρ .
 - 1. The best estimate of ρ is ρ itself.
 - i.e., $f(\rho:\rho) > f(\rho:\sigma)$ for all $\sigma \neq \rho$.
 - 2. $f(\rho : \sigma)$ should correspond to an operational test
 - i.e., someone's utility (reward or cost) for some practical procedure.

Quantum Strictly Proper Scoring Rules

Victor the Verifier measures a copy of ρ , and pays you $R_i(\sigma)$ if he gets outcome i.

1)
$$f(\rho:\sigma) = \sum_{i} p(i)R_{i}(\sigma)$$

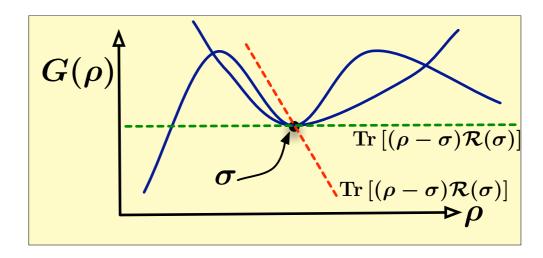
Measurement = $\{E_{i}\}: \sum_{i} E_{i} = 1$.

$$egin{aligned} f(
ho:\sigma) &= \sum_i ext{Tr}[
ho E_i] R_i(\sigma) \ &= ext{Tr}[
ho \mathcal{R}(\sigma)] \end{aligned}$$

where
$$\mathcal{R}(\sigma) = \sum_{i} E_{i} R_{i}(\sigma)$$
.

Define "value":
$$G(\rho) \equiv f(\rho : \rho) = \text{Tr}[\rho \mathcal{R}(\rho)]$$
.

2)
$$f(\rho:\rho) > f(\rho:\sigma)$$
 if $\sigma \neq \rho$.
 $f(\rho:\rho) > f(\sigma:\sigma) + f(\rho:\sigma) - f(\sigma:\sigma)$
 $G(\rho) > G(\sigma) + \text{Tr} [(\rho - \sigma)\mathcal{R}(\sigma)]$



- ullet $f(
 ho:\sigma)$ is a subtangent to G(
 ho)
- $G(\rho)$ must be strictly convex.
- The estimator's expected loss for lying is $\Delta(\rho:\sigma) \equiv G(\rho) f(\rho:\sigma)$.
- \bullet $\Delta(\rho:\sigma)$ is an operational divergence = good measure of σ 's predictive accuracy.
- \bullet Operational divergences are 1:1 with strictly convex "entropies" $G(\rho)$.

A survey of metrics

- Good Metrics (Operational Divergences)
 - 1. L2 distance: $\text{Tr}\left[(\rho-\sigma)^2\right] \Leftrightarrow R_i = 2s_i \text{Tr}(\sigma^2) 1$
 - 2. Relative entropy: $S(\rho||\sigma) = \text{Tr} \left|\rho\log\frac{\sigma}{\rho}\right| \Leftrightarrow R_i = \ln(s_i)$
- BAD METRICS
 - 1. Overlap: $1 \text{Tr}[\rho\sigma]$ Improper (motivates lying)

 - 2. Trace-norm: $||\rho \sigma||_1$ 3. Fidelity: $1 \left[\text{Tr}\sqrt{\sqrt{\sigma}\rho\sqrt{\sigma}}\right]^2$ \nearrow \nearrow

Bayesian Mean Estimation

unconditionally optimizes every operational divergence

If you know ρ , the optimal estimate is $\sigma = \rho$.

But what if your knowledge is uncertain (probabilistic)?

 \Longrightarrow the state could be ρ_i with probability π_i .

Key fact: $f(\rho : \sigma) = \text{Tr}[\rho \mathcal{R}(\sigma)]$ is linear in ρ .

 $\implies expected \text{ utility is}$

$$\overline{f} = \sum \pi_i f(
ho_i:\sigma) = f\left(\sum \pi_i
ho_i:\sigma
ight)$$

 \implies optimal estimate is $\sigma = \overline{\rho} \equiv \sum \pi_i \rho_i$

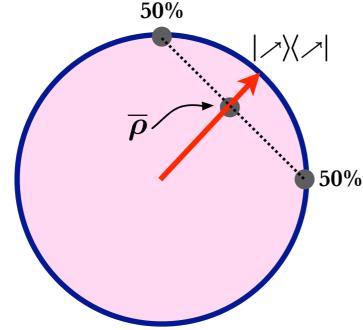
This is fairly straightforward, but tedious... see quant-ph/0603116

If unknown ρ was selected from distribution $\pi_0(\rho)d\rho$, and measurements $\mathcal{M} = \{E_1, E_2 \ldots\}$ were made, then:

- 1. Your knowledge is $\pi(\rho) d\rho = \frac{p(\mathcal{M}|\rho)\pi_0(\rho)d\rho}{\int p(\mathcal{M}|\rho)\pi_0(\rho)d\rho}$.
- 2. The optimal estimate is $\overline{\rho} = \int \rho \pi(\rho) d\rho$.

But isn't this obvious? NO

(a) Suppose we optimize fidelity.



$$\overline{F}\left(
ho_i, \overline{
ho}
ight) = rac{3}{4} \ \overline{F}\left(
ho_i, |
ho
angle
angle /
ho|) = rac{3+2\sqrt{2}}{4+2\sqrt{2}} pprox 0.85$$

MORAL: fidelity measures how well σ simulates ρ , not how well it estimates ρ .

- (b) Suppose we assume the future will look (statistically) just like the past.
- We know what future datasets will look like
- We can add them to measurement record
- $\bullet \ \mathcal{M} \longrightarrow \mathcal{M} \cup \mathcal{M} \cup \mathcal{M} \cup \ldots$
- $ullet p(\mathcal{M}|
 ho) \longrightarrow p(\mathcal{M}|
 ho)^{\infty}$
- So $\pi(\rho) d\rho \rightarrow \delta (\rho \rho_{\text{MLE}})$

MORAL: MLE can be derived by assuming this "frequentist axiom".

(This explains a lot...)

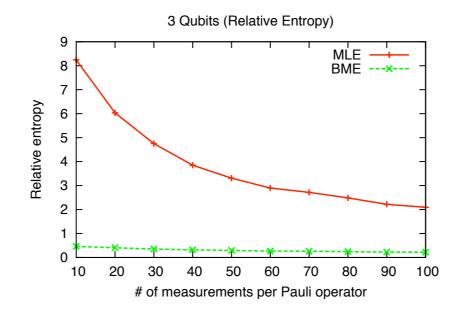
How good is MLE?

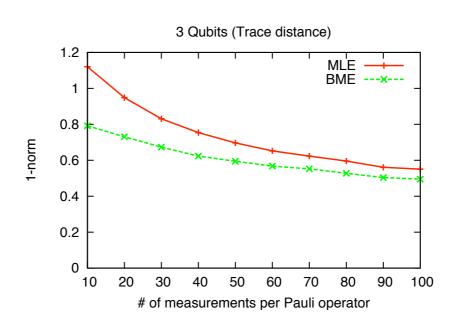
NUMERICS:

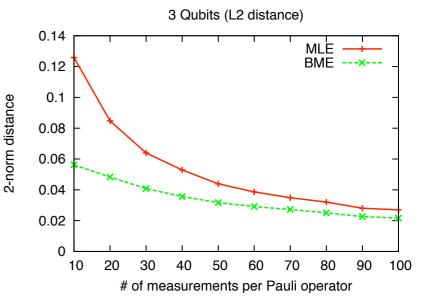
- 1. Generate random (Hilbert-Schmidt measure) mixed states ρ for n qubits,
- 2. Measure each of 4^{n-1} Pauli observables N times,
- 3. Analyze measurement record to get σ ,
- 4. Compare σ to ρ using:
 - (a) L2-norm, $\operatorname{Tr}[(\rho \sigma)^2],$
 - (b) relative entropy, $\text{Tr} \left[\rho (\ln \rho \ln \sigma) \right],$
 - (c) fidelity,

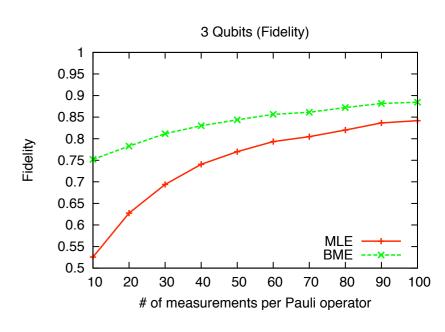
$$\operatorname{Tr}\left[\sqrt{\sqrt{\sigma}
ho\sqrt{\sigma}}
ight]^2$$

(d) L1-norm, $\operatorname{Tr}[|\boldsymbol{\rho} - \boldsymbol{\sigma}|].$





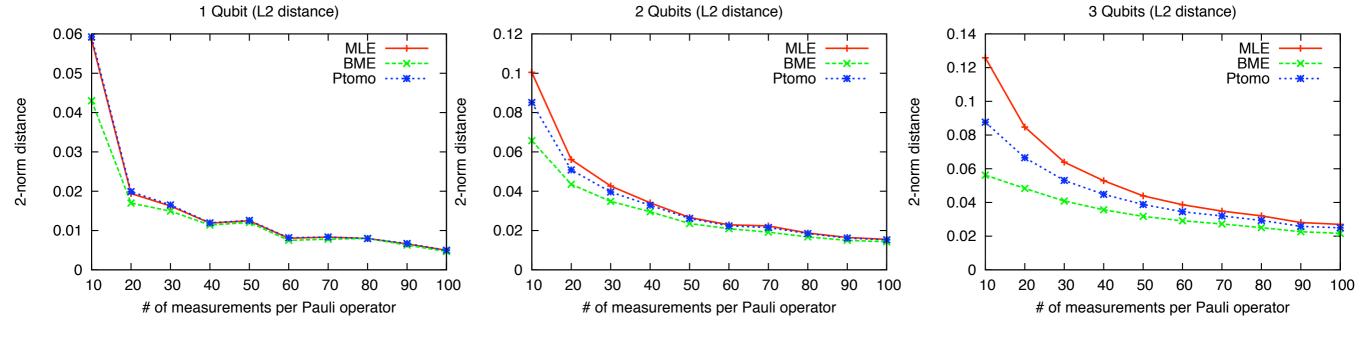




Quick & Dirty Tomography

- 1. MLE gets computationally difficult for large systems.
- 2. BME is even harder!
- 3. $\frac{1}{\sqrt{3}}$ (|D.James| + |T.Havel| + |P.Jessen|) suggested "quick 'n' dirty tomography".

$$\sigma_{
m tomo} = egin{pmatrix} \lambda_1 & & & & \\ \lambda_2 & & & \\ -\lambda_3 & & -\lambda_4 \end{pmatrix} \longrightarrow egin{pmatrix} \lambda_1 & & & \\ \lambda_2 & & & \\ & 0 & & \\ & & 0 \end{pmatrix} \longrightarrow rac{1}{\lambda_1 + \lambda_2} egin{pmatrix} \lambda_1 & & & \\ \lambda_2 & & & \\ & & 0 & \\ & & 0 \end{pmatrix} = \sigma$$



Conclusions

- It's worth thinking carefully and deeply about state estimation.
- Operational divergences are a good way to evaluate the predictive accuracy of estimates.
- BME is optimal, and a baseline for evaluating other (more efficient) approaches.
- Quick & dirty methods can outperform MLE.