

Quantum gates in Kane's model based on adiabatic controlling processes and verification of its physical realization

July, 19–21, 2004

Quantum Information and Quantum Control Conference
the University of Toronto and the Fields Institute

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Motivation & Aim

- Physical realization of the Kane's model (Nature 393, 133 (1998))
 - difficult, but challenging problem !!
 - possibility: **Single Ion Injection Method (SII)**
(T. Shinada, *et al.*, Jpn. Appl. Phys. 41, L287 (2002)
(Waseda University))
- Control of quantum systems by classical manipulation
 - Kane's model: change the gate voltage
⇒ change **the strength of interaction (locally)**

We discuss **the construction of quantum gates**,
based on **rigorous analyses in the proposed model.**

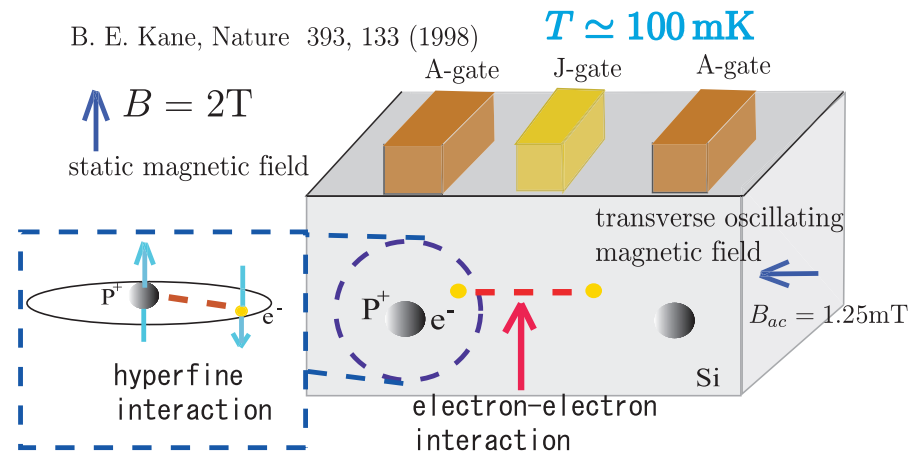
Model

N-qubit system

$$H = H_0 + H_{ac}, \quad H_0 = \sum_{i=1}^N H^i + \sum_{\langle i,j \rangle} J_{ij} \sigma^{ie} \cdot \sigma^{je}, \quad H_{ac} = \sum_{i=1}^N H_{ac}^i$$

$$H^i = \mu_B B \sigma_z^{ie} - g_n \mu_n B \sigma_z^{in} + A_i \sigma^{ie} \cdot \sigma^{in}, \quad H_{ac}^i = B_{ac} \mathbf{m}^i \cdot (\mu_B \sigma^{ie} - g_n \mu_n \sigma^{in})$$

- A_i : hyperfine interaction (HF) (on the i -th dopant)
- J_{ij} : electron-electron exchange interaction (EE) (between i -th and j -th dopant), $j = i+1$
- (gate voltage)=0 $\Rightarrow A_i = A_0, J_{ij} = 0$
($2A_0/h = 58$ MHz)
- $\mu_B B = 1.158$ meV,
 $g_n \mu_n B = 7.135 \times 10^{-5}$ meV
- $\mathbf{m}^i = (\cos(\omega_{act}), -\sin(\omega_{act}), 0)$



Method – adiabatic controlling processes –

- A_i, J_{ij} : **adiabatic controlling processes**
- **transverse magnetic field: instantaneously switch on–off** ($\because B_{ac}/B$: small)

$$2A_0/\mu_B B \simeq 2.07 \times 10^{-4}, g_n \mu_n B/\mu_B B \simeq 0.62 \times 10^{-4}, J \sim \mu_B B, B_{ac}/B \simeq 10^{-3}$$

- ★ **phase shift for i -th qubit**
- ★ **spin flip for i -th qubit**
- ★ **controlled-Z between i -th and j -th qubits ($j = i + 1$)**

Diagonalization of Hamiltonian

- H^i : related with the dynamics for one qubit

$$S^i = (\sigma_z^{ie} + \sigma_z^{in})/2, [S^i, H^i] = 0$$

$$H^i = E_{\uparrow 0}^i |u_{\uparrow 0}^i\rangle \langle u_{\uparrow 0}^i| + E_{\uparrow 1}^i |u_{\uparrow 1}^i\rangle \langle u_{\uparrow 1}^i| + E_{\downarrow 0}^i |u_{\downarrow 0}^i\rangle \langle u_{\downarrow 0}^i| + E_{\downarrow 1}^i |u_{\downarrow 1}^i\rangle \langle u_{\downarrow 1}^i|$$

$$|u_{\downarrow 0}^i\rangle = (-2A_i |\uparrow 1\rangle + (\epsilon + \sqrt{\epsilon^2 + 4A_i^2}) |\downarrow 0\rangle) / N_i, |u_{\downarrow 1}^i\rangle = |\downarrow 1\rangle$$

$$E_{\uparrow 0}^i = \epsilon - 2g_n \mu_n B + A_i, E_{\uparrow 1}^i = -A_i + \sqrt{\epsilon^2 + 4A_i^2},$$

$$E_{\downarrow 0}^i = -A_i - \sqrt{\epsilon^2 + 4A_i^2}, E_{\downarrow 1}^i = -\epsilon + 2g_n \mu_n B + A_i \quad (\epsilon = \mu_B B + g_n \mu_n B)$$

- $H^{ij} = H^i + H^j + J_{ij} \sigma^{ie} \cdot \sigma^{je}, j = i + 1$

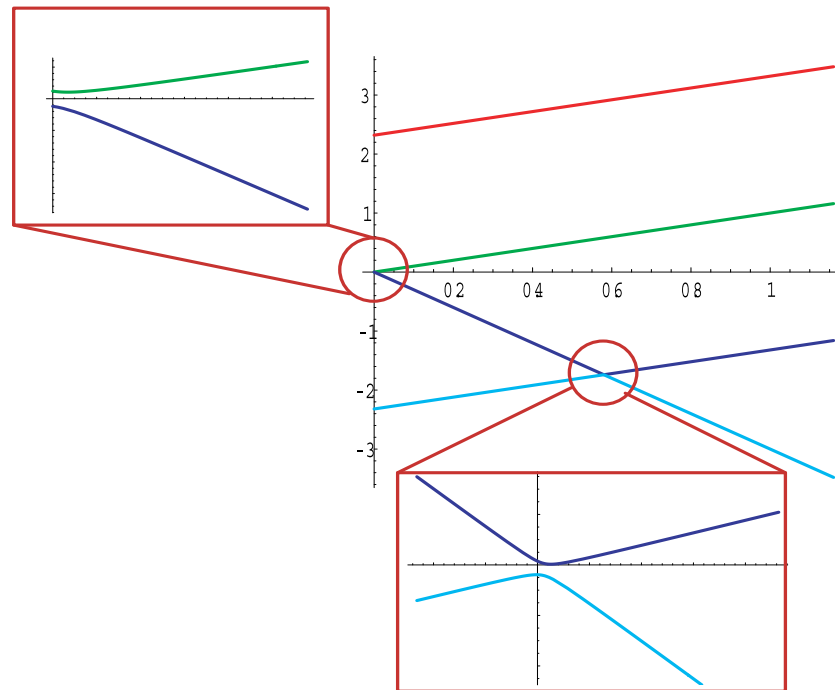
$$S^{ij} = S^i + S^j, [S^{ij}, H^{ij}] = 0, P^{ij}: \text{exchange of labels for identical particles}$$

$$A_i = A_j \iff [P^{ij}, H^{ij}] = 0$$

$\Rightarrow H^{ij}$: block diagonal form \Rightarrow **analytical diagonalization**

quantum number: $S^{ij} \rightarrow s = 2, 1, 0, -1, -2, P^{ij} \rightarrow p = +, -$

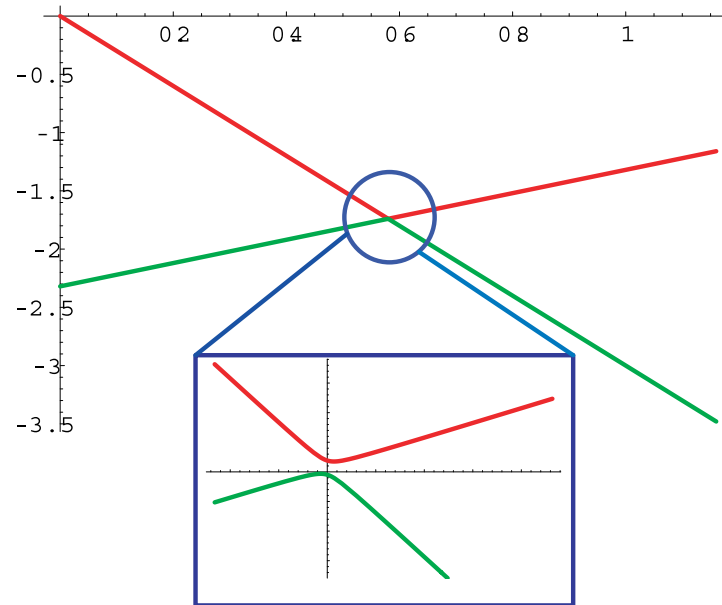
Energy level I



vertical axis: eigenvalues [meV] for $(s, p) = (0, +)$, horizontal axis: J [meV],

cyan line: $|u_{\downarrow 0}^i\rangle|u_{\downarrow 0}^{i+1}\rangle$ ($J = 0$)

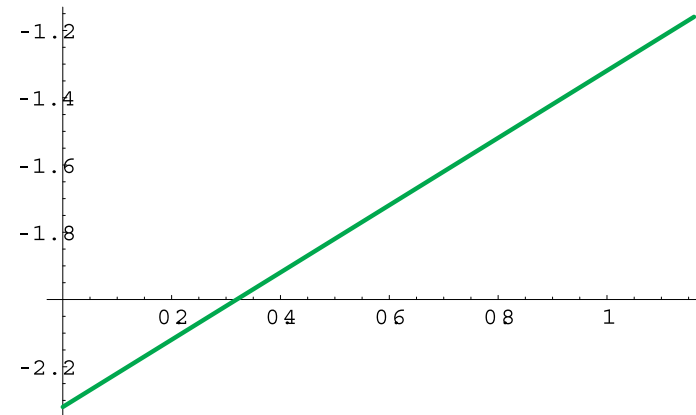
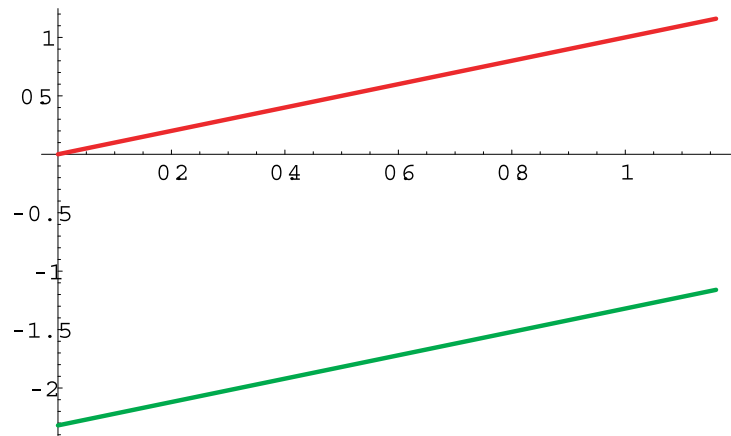
Energy level II



vertical axis: eigenvalues [meV] for $(s, p) = (-1, -)$, horizontal axis J [meV],

green line: $(|u_{\downarrow 0}^i\rangle|u_{\downarrow 1}^{i+1}\rangle - |u_{\downarrow 1}^i\rangle|u_{\downarrow 0}^{i+1}\rangle)/\sqrt{2}(J = 0)$

Energy level III



- left figure, vertical axis: eigenvalues [meV] for $(s, p) = (-1, +)$, horizontal axis: J [meV],
green line: $(|u_{\downarrow 0}^i\rangle|u_{\downarrow 1}^{i+1}\rangle + |u_{\downarrow 1}^i\rangle|u_{\downarrow 0}^{i+1}\rangle)/\sqrt{2}$ ($J = 0$)
- right figure, vertical axis: eigenvalue [meV] for $(s, p) = (-2, +)$, horizontal axis: J [meV],
green line: $|u_{\downarrow 1}^i\rangle|u_{\downarrow 1}^{i+1}\rangle$ ($J = 0$)

Representation of Quantum Information

Eigenvector for H^i : $|u_{\uparrow 0}^i\rangle (= |\downarrow 0\rangle)$, $|u_{\uparrow 1}^i\rangle (\simeq |\uparrow 1\rangle)$, $|u_{\downarrow 0}^i\rangle (\simeq |\downarrow 0\rangle)$, $|u_{\downarrow 1}^i\rangle (= |\downarrow 1\rangle)$

- **Initialization:** $T \simeq 100 \text{ mK} \Rightarrow \rho = \frac{1}{Z} \exp\left(-\beta \sum_{i=1}^N H^i\right) \simeq \bigotimes_{i=1}^N |u_{\downarrow 0}\rangle_i \langle u_{\downarrow 0}|$
- $|E_{\downarrow 0}^i - E_{\downarrow 1}^i|$: characterized by A_0
 \Rightarrow **the gate construction through the control of HF**

$$\boxed{|u_{\downarrow 0}^i\rangle \doteq |0\rangle_L, |u_{\downarrow 1}^i\rangle \doteq |1\rangle_L}$$

- **controlled operation (controlled-Z) ($j = 1 + 1$)**

computational basis: $|u_{\downarrow 0}^i\rangle|u_{\downarrow 0}^j\rangle, |u_{\downarrow 0}^i\rangle|u_{\downarrow 1}^j\rangle, |u_{\downarrow 1}^i\rangle|u_{\downarrow 0}^j\rangle, |u_{\downarrow 1}^i\rangle|u_{\downarrow 1}^j\rangle$

$$|v_1\rangle \equiv |u_{\downarrow 0}^i\rangle|u_{\downarrow 0}^j\rangle \leftrightarrow (s, p) = (0, +), |v_4\rangle \equiv |u_{\downarrow 1}^i\rangle|u_{\downarrow 1}^j\rangle \leftrightarrow (s, p) = (-2, +),$$

$$|v_{\pm}\rangle \equiv (|u_{\downarrow 0}^i\rangle|u_{\downarrow 1}^j\rangle \pm |u_{\downarrow 1}^i\rangle|u_{\downarrow 0}^j\rangle)/\sqrt{2} \leftrightarrow (s, p) = (-1, \pm)$$

Each vector belongs to the different subspace !

Result 1 – phase shift –

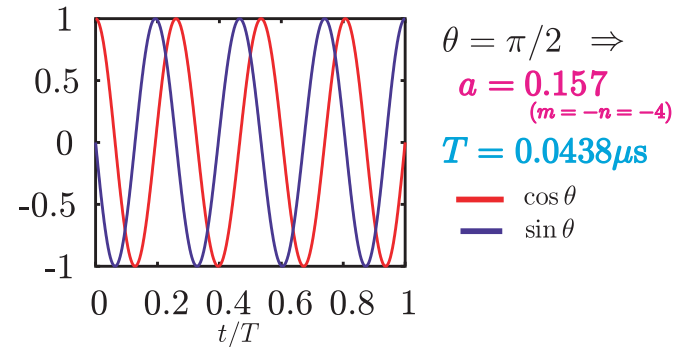
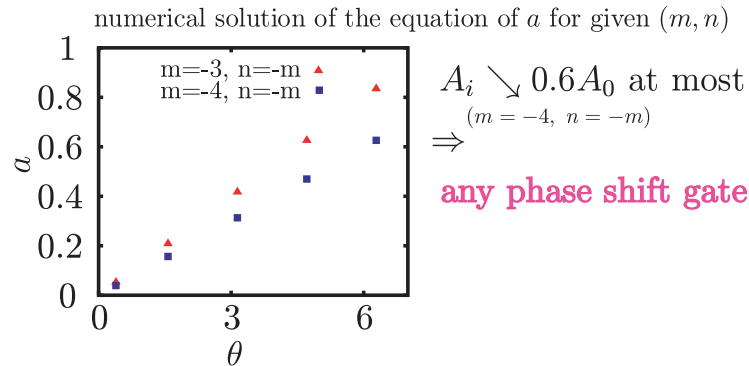
$J_{kl} = 0$, $A_i(t) = A_0(1 - a \sin(\pi t/T_{op}))$, $A_k = A_0$ ($k \neq i$), T_{op} : operation time, a : parameter

- phase difference between $|u_{\downarrow 0}^i\rangle$ (**adiabatic**) and $|u_{\downarrow 1}^i\rangle$ (**eigenstate**)

i th qubit: $\Theta_i = \frac{T_{op}}{\hbar} \left(-2 \int_0^1 A_i(\tau) d\tau - \int_0^1 \sqrt{\epsilon^2 + 4A_i(\tau)^2} d\tau + \epsilon - 2g_n \mu_n B \right)$

others: $\Theta_0 = \frac{T_{op}}{\hbar} \left(-2A_0 - \sqrt{\epsilon^2 + 4A_0^2} + \epsilon - 2g_n \mu_n \right)$

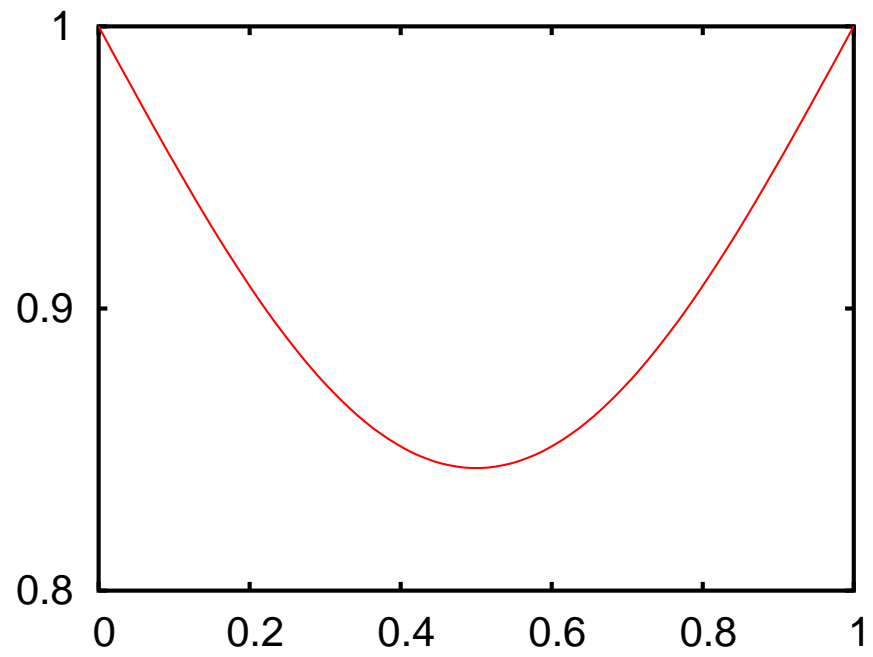
$\Rightarrow \theta = \Theta_i - \Theta_0 - 2n\pi$, ($\Theta_0 = 2m\pi$, $m \in \mathbb{Z}$, $n \in \mathbb{Z}$) \rightarrow determine the value of a



$|u_{\downarrow 1}\rangle$: **eigenstate**, $|\langle u_{\downarrow 0} | \psi(T_{op}) \rangle| = 1 + C/\epsilon T_{op}$, $|\langle u_{\uparrow 1} | \psi(T_{op}) \rangle| = C/\epsilon T_{op}$, $C \simeq 10^{-5}$,

$\epsilon T_{op} \simeq 10^3 \Rightarrow$ **adiabatic approximation: good** \Rightarrow **error** $\sim 10^{-8}$

profile of A_i



vertical axis: A_i/A_0 , horizontal axis: t/T_{op}

Result 2 – spin sfip –

Larmor resonance frequency

$$\hbar\omega_{ac} = -\epsilon + 2g_n\mu_n B + 2A + \sqrt{\epsilon^2 + 4A^2}$$

$A_i = A$, $A_j = A_0 (j \neq i)$: control the resonance condition locally

operation time $T_{op} = \hbar\alpha/\nu_\theta B_{ac}$, $\nu_\theta \simeq g_n\mu_n$, α : angle (i.e., $e^{-i\alpha\sigma_x}$)

This method correspond to Hill & Goan's work PRA 012321 (2003)

★ Hamiltonian in the rotating frame

$\Rightarrow \exists$ the term related with transition between the different electron spin states

$$\text{error: } \|\psi_{rot}(t) - \psi_d(t)\| \leq \left(2 + \frac{\mu_{-\theta}^i \alpha}{\nu_\theta^i}\right) \frac{2\mu_{-\theta}^i B_{ac}}{\epsilon - g_n\mu_n B_{ac} \cos \theta^i}$$

$$\mu_\theta^i = \mu_B \cos \theta^i - g_n\mu_n \sin \theta^i, \quad \nu_\theta^i = \mu_B \sin \theta^i + g_n\mu_n \cos \theta^i \quad \cos \theta^i = (\epsilon + \sqrt{\epsilon^2 + 4A_i^2})/N_i, \quad \sin \theta^i = 2A_i/N_i$$

$B_{ac}/B \sim 10^{-3}$: often used \Rightarrow error \rightarrow large !!

Result 3 – controlled-Z –

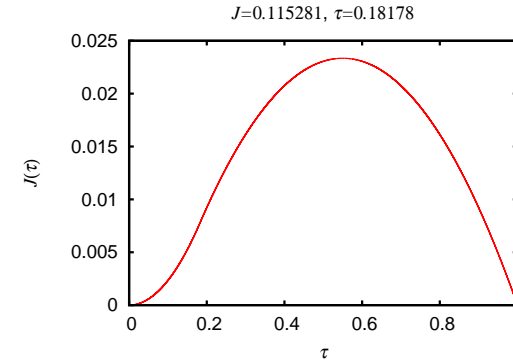
- For example

$$J(\tau) = \begin{cases} J_c \alpha \tau^2 & 0 \leq \tau < \tau_c \\ J_c \{\beta - (\tau - \tau_0)^2\} & \tau_c < \tau \leq 1 \end{cases}$$

$$\tau = t/T_{op}, \quad \alpha = (1 - \tau_c)^2 / \{\tau_c(1 - \tau_c)\},$$

$$\beta = (1 - \tau_c)^2 / (2 - \tau_c)^2, \quad \tau_0 = 1 / (2 - \tau_c)$$

parameter: τ_c , J_c , and T_{op}



adiabatic time evolution \Leftarrow analytical calculation !

$$|v_1\rangle \rightarrow e^{i\delta_1} |v_1\rangle, \quad |v_{\pm}\rangle \rightarrow e^{i\delta_{\pm}} |v_{\pm}\rangle = e^{i\delta_1} e^{i(\delta_{\pm} - \delta_1)} |v_{\pm}\rangle, \quad |v_4\rangle \rightarrow e^{i\delta_4} |v_4\rangle = e^{i\delta_1} e^{i(\delta_4 - \delta_1)} |v_4\rangle$$

controlled-Z $\iff \delta_{\pm} - \delta_1 = 2m_{\pm,1}\pi, \delta_4 - \delta_1 = 2m_{4,1}\pi + \pi, m_{\pm,1}, m_{4,1} \in \mathbb{Z}$

J_c/ϵ	τ_c	$m_{4,1}$	$m_{+,1}$	$m_{-,1}$	$T_{op} [\mu s]$
0.115281	0.181788	25	25	-24	0.0054
0.695156	0.0575511	50	50	-49	0.0054

adiabatic approximation: good (numerically)

Matrix for the controlled-Z

$$e^{i\delta_1} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & a_+ & a_- & 0 \\ 0 & a_- & a_+ & 0 \\ 0 & 0 & 0 & a \end{pmatrix} \begin{matrix} |00\rangle \\ |01\rangle \\ |10\rangle \\ |11\rangle \end{matrix}$$

$$a_+ = (e^{i(\delta_+ - \delta_1)} + e^{i(\delta_- - \delta_1)})/2 = (e^{i2m_+,1\pi} + e^{i2m_-,1\pi})/2 = 1$$

$$a_- = (e^{i(\delta_+ - \delta_1)} - e^{i(\delta_- - \delta_1)})/2 = (e^{i2m_+,1\pi} - e^{i2m_-,1\pi})/2 = 0$$

$$a = e^{i(\delta_4 - \delta_1)} = e^{i2m_4,1\pi} e^{i\pi} = -1$$

⇓

$$e^{i\delta_1} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{matrix} |00\rangle \\ |01\rangle \\ |10\rangle \\ |11\rangle \end{matrix}$$

Conclusion

- We discuss the construction of the quantum gates through the adiabatic controlling processes
 - phase shift gate: we determine the good parameter.
adiabatic approximation: **good** \rightarrow very small error
 - spin flip gate: standard value ($B_{ac}/B \sim 10^{-3}$) \Rightarrow very large error
improvement: **the value of $B_{ac}/B \Rightarrow$ smaller**
 \Rightarrow But, **the operation time increases** $\Rightarrow \exists$ Optimal value ?
 - controlled-Z gate: we show a possible several sets of parameters.

Examine much more.

perspective

- The theoretical estimation of the required accuracy of “ion injection” which enables us to perform the quantum computation.